

# DESIGNING ROBUST QUANTUM REFRIGERATORS IN DISORDERED SPIN MODEL

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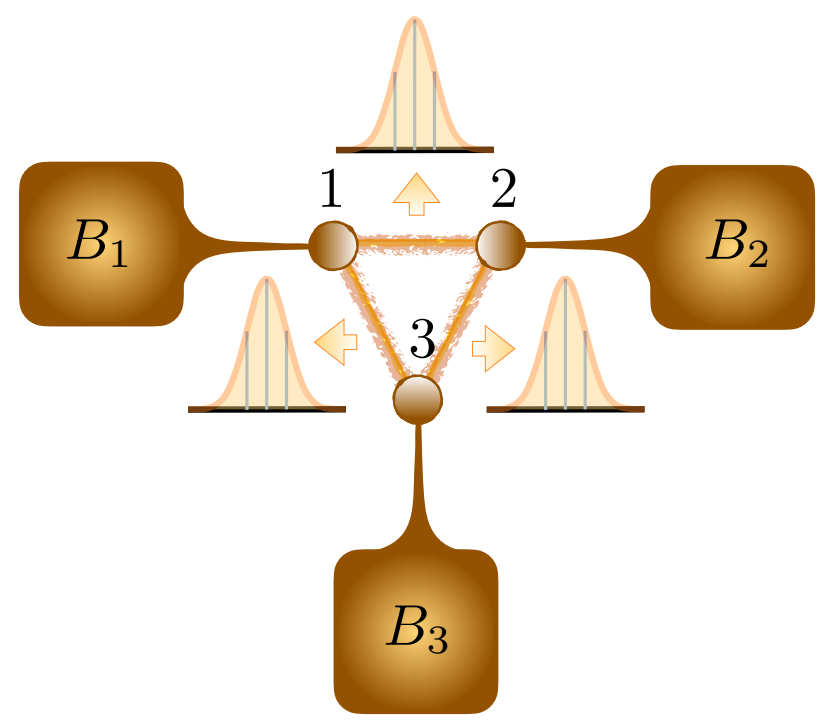
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## Abstract

We explore a small quantum refrigerator in which the working substance is made of paradigmatic nearest-neighbor quantum spin models, the  $XYZ$  and the  $XY$  model with Dzyaloshinskii-Moriya interactions, consisting of two and three spins, each of which is in contact with a bosonic bath. We identify a specific range of interaction strengths which can be tuned appropriately to ensure a cooling of the selected spin in terms of its local temperature in the weak-coupling limit. Moreover, we report that in this domain, when one of the interaction strengths is disordered, the performance of the thermal machine operating as a refrigerator remains *almost* unchanged instead of degradation, thereby establishing the flexibility of this device. However, to obtain a significant amount of cooling via ordered as well as disordered spin models, we observe that one has to go beyond the weak-coupling limit and compute the figures of merits by using global master equations.

## Design of Refrigerator



$$H_F = \sum_{i=1}^N h_i \sigma_z^i, \quad (1)$$

$$H_{xy} = \sum_{i=1}^N J_{i,i+1}^{xy} [(1+\gamma)\sigma_x^i \sigma_x^{i+1} + (1-\gamma)\sigma_y^i \sigma_y^{i+1}] \quad (2)$$

$$H_z = \sum_{i=1}^N J_{i,i+1}^z \sigma_z^i \sigma_z^{i+1}, \quad (3)$$

$$H_{dm} = \sum_{i=1}^N J_{i,i+1}^{dm} (\sigma_x^i \sigma_y^{i+1} - \sigma_y^i \sigma_x^{i+1}). \quad (4)$$

- $\gamma$  is the  $xy$  anisotropy parameter,  $h_i$  is the strength of the local magnetic field acting on the spin  $i$
- The initial state of the system is given by  $\rho_s^0 = \bigotimes_{i=1}^N \rho_i^0$ , where  $\rho_i^0 = \exp(-\beta_i^0 h_i \sigma_z^i) / \text{Tr}[\exp(-\beta_i^0 h_i \sigma_z^i)]$ , with  $\beta_i^0 = (k_B T_i^0)^{-1}$ ,  $k_B$  is the Boltzmann constant.

## Master Equations

- Local Master Equation

$$\mathcal{D}_i(\rho) = \Gamma_i \left[ (n_\omega^i + 1) (\sigma_i^- \rho \sigma_i^+ - \frac{1}{2} \{ \sigma_i^+ \sigma_i^-, \rho \}) + n_\omega^i (\sigma_i^+ \rho \sigma_i^- - \frac{1}{2} \{ \sigma_i^- \sigma_i^+, \rho \}) \right], \quad (5)$$

- Global Master Equation

$$\mathcal{D}(\rho) = \sum_{\omega>0} \gamma_\omega^i \left[ (A_\omega^i \rho A_\omega^{i\dagger} - \frac{1}{2} \{ A_\omega^{i\dagger} A_\omega^i, \rho \}) + (A_\omega^{i\dagger} \rho A_\omega^i - \frac{1}{2} \{ A_\omega^i A_\omega^{i\dagger}, \rho \}) \right], \quad (6)$$

where the operator  $A_\omega^i$  is given by

$$e^{iH_s t} (\sigma_i^+ + \sigma_i^-) e^{-iH_s t} = 2 \sum_{\omega} A_\omega^i e^{-i\omega t} \quad (7)$$

## Local Refrigeration

- Local Temperature

$$T_i(t) = \frac{2h_i}{\ln[\tau_i(t)^{-1} - 1]} \quad (8)$$

- A local steady-state cooling of the spin  $i$  is achieved if

$$T_i^s = T_i(t \rightarrow \infty) < T_i^0 \quad (9)$$

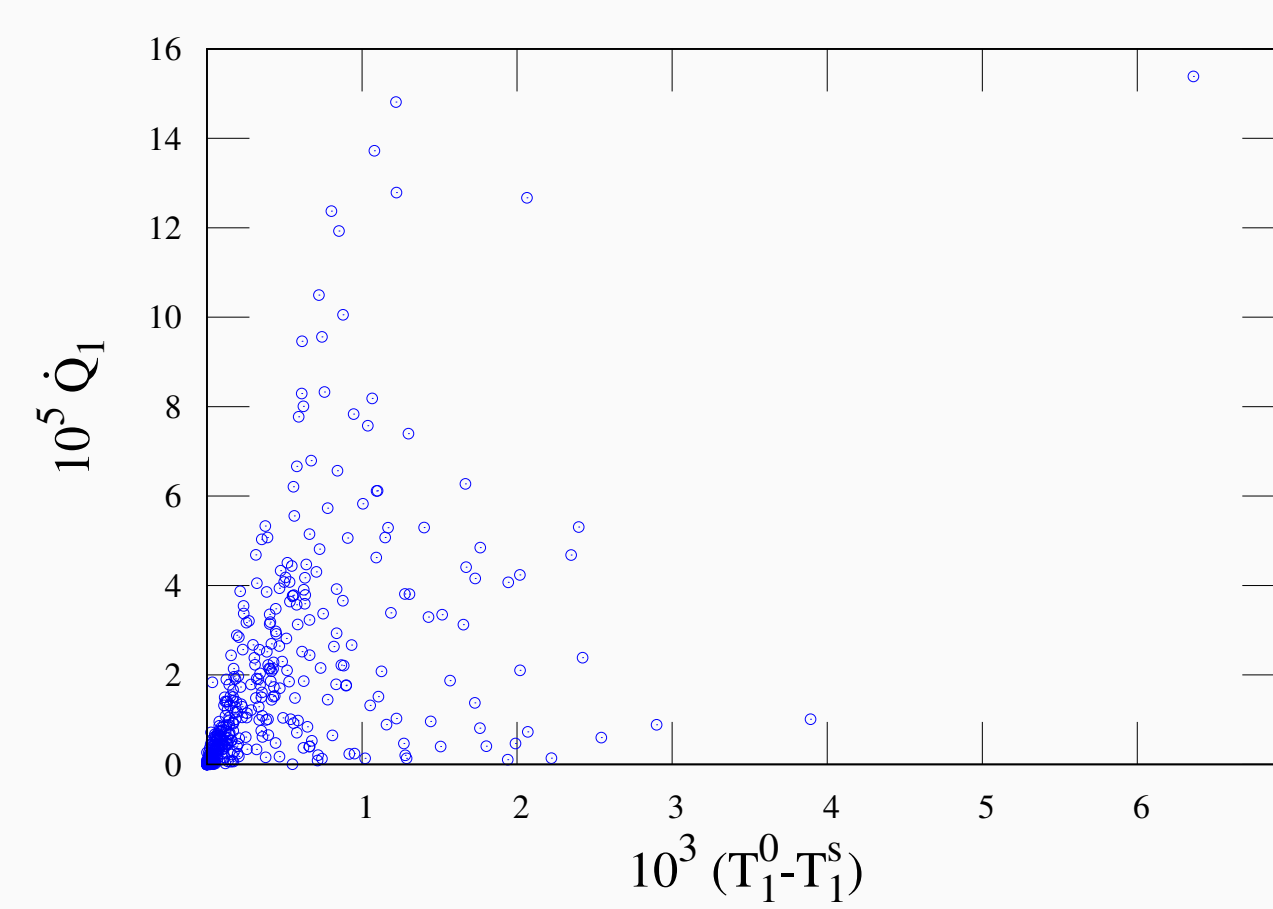
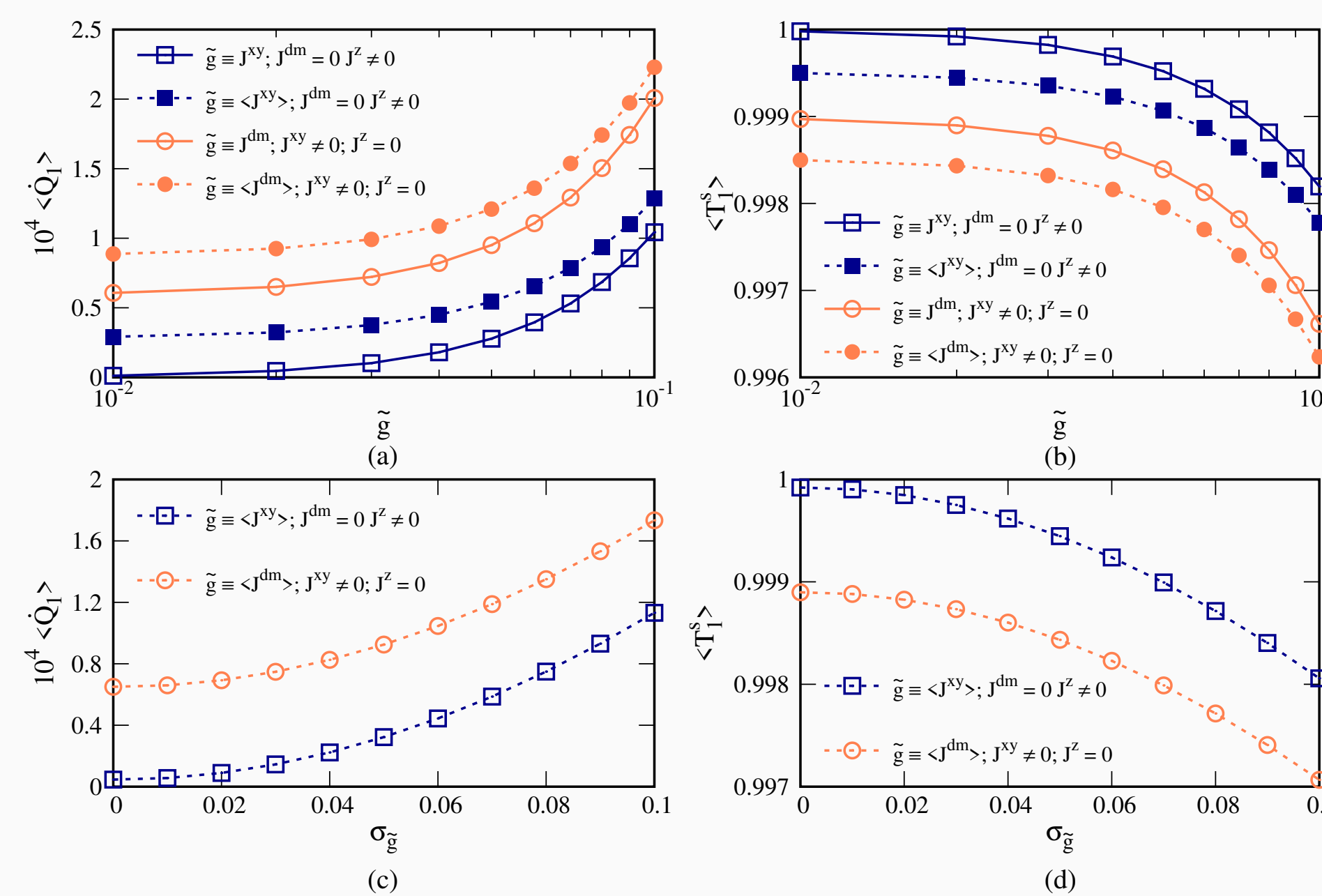
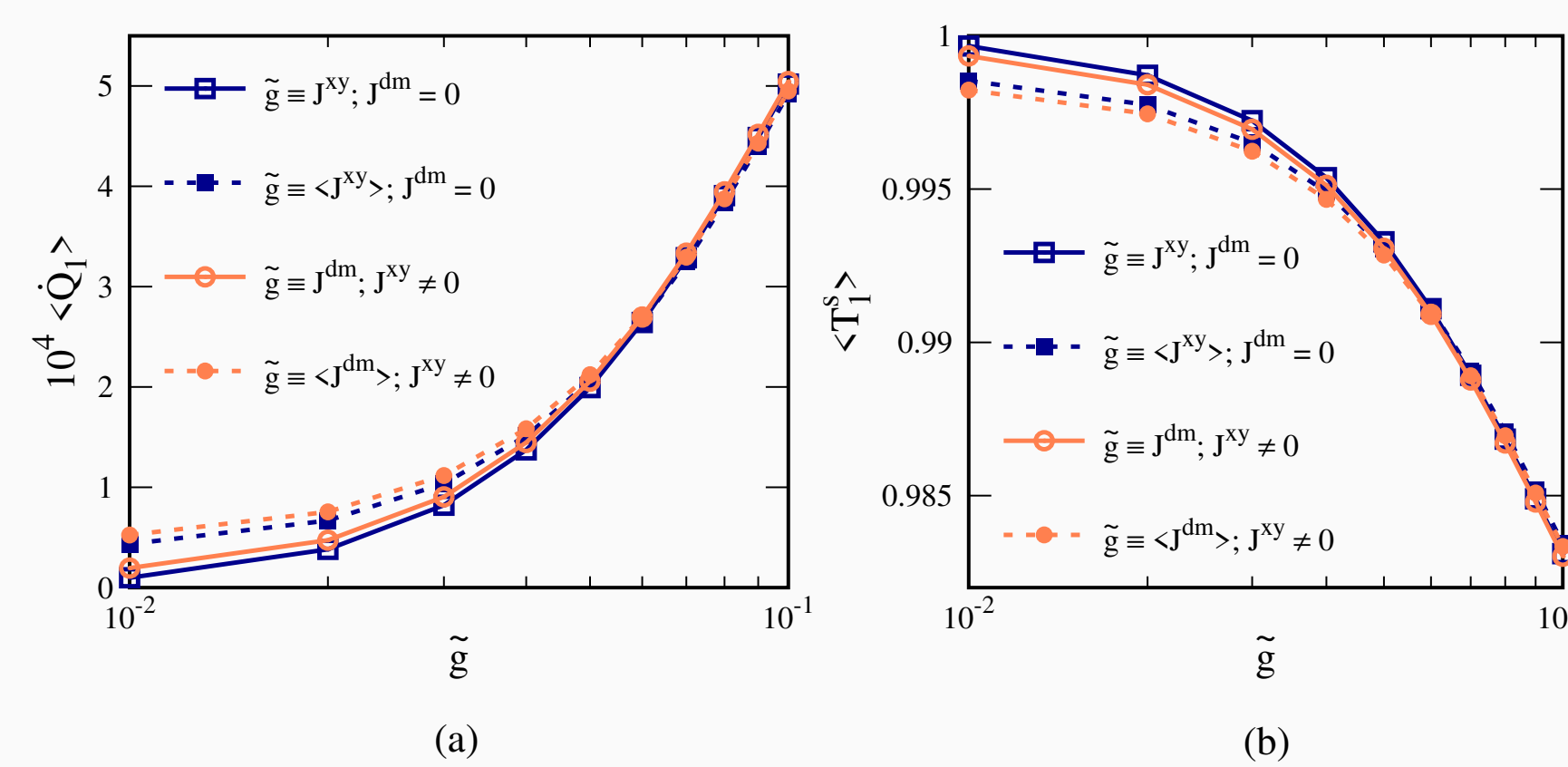
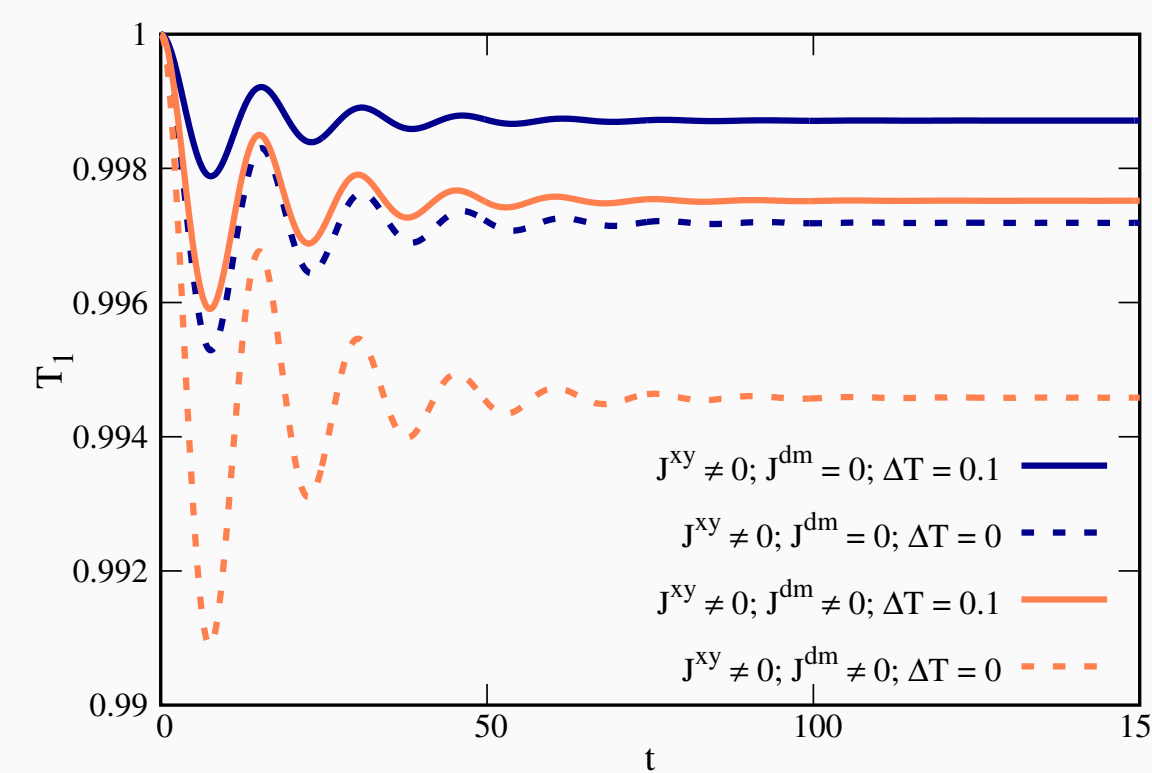
where  $T_i^0$  denotes initial temperature of the spin  $i$

## Disorder

$$\langle \mathcal{Q}(\langle g \rangle, \sigma_g) \rangle = \int \mathcal{P}(g) \mathcal{Q}(g) d(g), \quad (10)$$

where  $g$  is the parameter values, which are chosen from a Gaussian distribution ( $\mathcal{P}(g)$ ) of mean  $\langle g \rangle$  and standard deviation  $\sigma_g$  quantifying the strength of the disorder.

## Quantum refrigerator: Order vs. disorder



- A non-zero  $XY$  interaction strength,  $J^{xy}$ , results in an evolution of the system, leading to a local cooling of spin 1, irrespective of the value of  $J^z$ .

- More importantly, we report that vanishing  $\Delta T = T_2^0 - T_1^0$  proves to be advantageous with respect to cooling than that of a non-vanishing  $\Delta T$  if we suitably adjust the parameters of  $H_s$  and the spin-bath interaction strength (comparing solid ( $T_2^0 = 1.1, T_1^0 = 1$ ) and dashed lines ( $T_2^0 = T_1^0 = 1$ )).

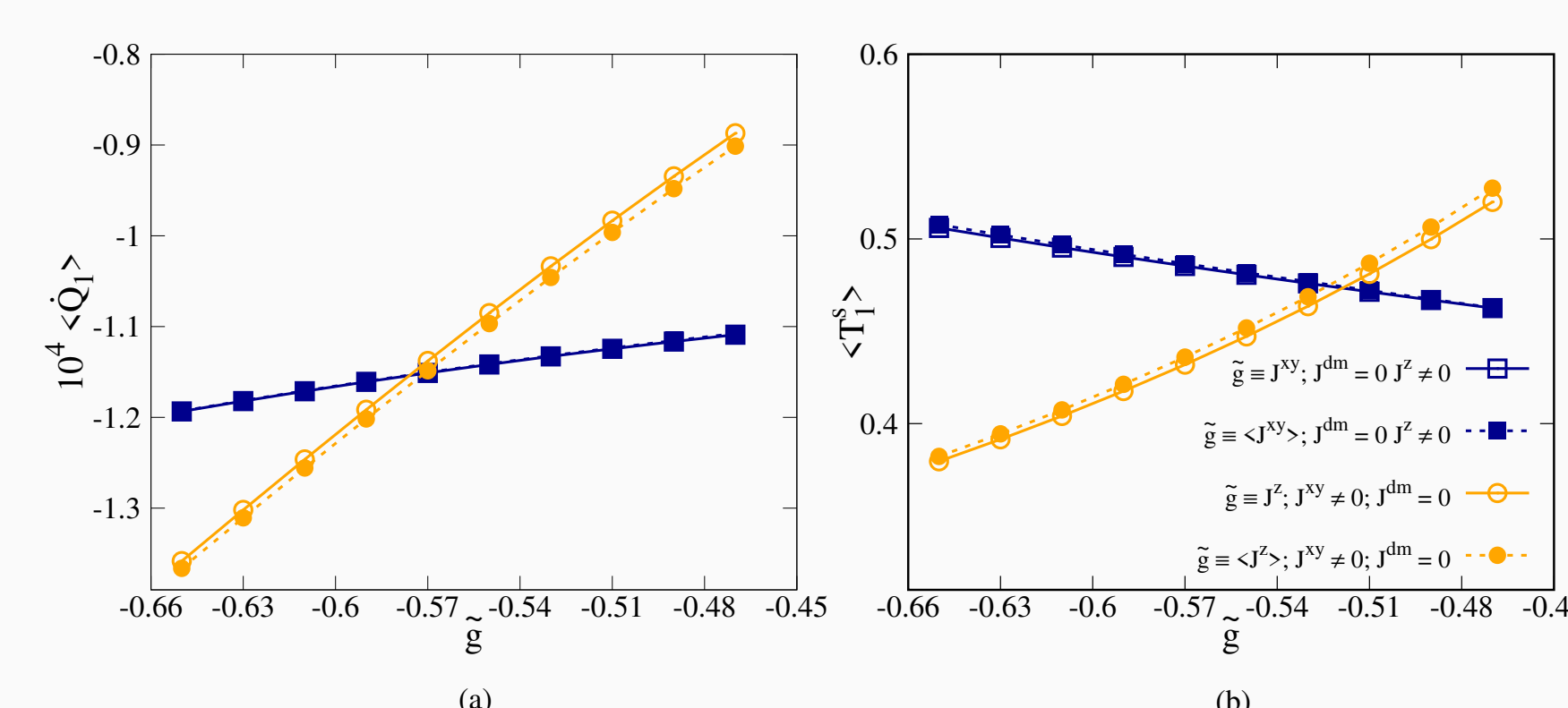
The heat current (the steady-state temperature) remains almost constant when the strength of the spin-exchange interaction is  $\leq 10^{-2}$ , and increases with an increase in the value of  $J^{xy}$  within the weak-coupling limit ( $\leq 10^{-1}$ ), irrespective of the presence of the interactions in the  $z$ -plane, i.e., independent of the values of  $J^z$ .

- We compare the performance of refrigerator designed via ordered and disordered three spin model. It shows robustness of refrigeration in presence of disorder in the system ((a)-(b)).

- We systematically increase the value of the disorder-strength up to  $10^{-1}$ , the average value of the heat current of the first spin attains a more positive value, while the steady-state temperature becomes lower ((c)-(d)) than that of the model with low disorder-strength. It clearly exhibits an advantage to attain a lower steady-state temperature of the refrigerated spin in the presence of disorder.

The values of the local magnetic fields,  $\{h_1, h_2, h_3\}$ , corresponding to the individual spins are chosen uniformly from  $[1.1, 5]$  while the values of the spin-bath interaction parameters  $\{\Gamma_1, \Gamma_2, \Gamma_3\}$  as well as the spin-exchange interaction strengths  $\{J^{xy}, J^z\}$  are chosen from a uniform distribution of range  $[0, 10^{-1}]$ . Here  $T_1^0 = 1, T_2^0 = 2$  and  $T_3^0 = 3$ . Among  $10^4$  choices of parameters, only 4.11% points are displayed for which local temperature of the first spin is lower than unity which indicate that there is no monotonic relation between them.

## Beyond weak-coupling limit



- We find that the steady-state temperature and the corresponding quenched averaged temperature of the first spin can substantially be decreased in the strong-coupling domain compared to that obtained in the weak-coupling limit.

- While the robustness of local cooling in the disordered refrigerator is a common feature in both local and global master equations, an interesting difference between these two situations emerge.

## References

1. N. Linden, S. Popescu, and P. Skrzypczyk, Phys. Rev. Lett. **105**, 130401 (2010).
2. T. K. Konar, S. Ghosh, A. K. Pal and A. Sen (De), Phys. Rev. A **105**, 022214 (2022).