

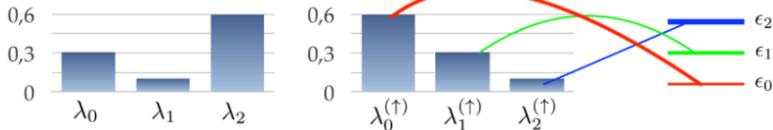
Optimal local work extraction from bipartite quantum systems in the presence of Hamiltonian couplings

Raffaele Salvia, Giacomo De Palma, Vittorio Giovannetti

Ergotropy

The *ergotropy* of a system is the **maximum work** that can be extracted from it with unitary operations

$$\mathcal{E}(\rho, H) = \text{Tr}[\rho H] - \min_U \text{Tr}[U \rho U^\dagger H]$$



It is a convex and non-linear functional of ρ and H

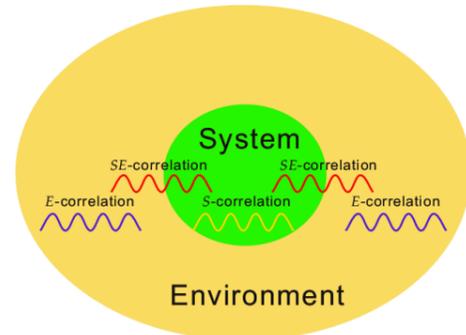
$$\mathcal{E}(\hat{\rho}) = 0 \text{ for passive states}$$

$$\mathcal{E}(\hat{\rho}) = E(\hat{\rho}) \text{ if and only if the state is pure}$$

Local ergotropy

We define the *local ergotropy* as the maximum extractable work with **local unitary operations**:

$$\mathcal{E}_S(\rho_{SE}, H_{SE}) = \text{Tr}[\rho_{SE} H_{SE}] - \min_{U_S} \text{Tr}[(U_S \otimes I_E) \rho_{SE} (U_S^\dagger \otimes I_E) H_{SE}]$$

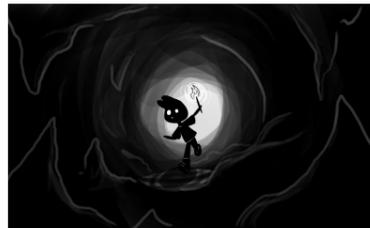
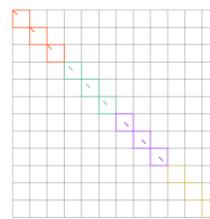


A difficult problem...

In the classical case (UNITARIES \rightarrow PERMUTATIONS), the problem reduces to finding, given a probability distribution p_{ij} over energies ϵ_{ij} ,

$$\min_{\pi} \sum_{i,j} p_{\pi(i),j} \epsilon_{i,j}$$

This is the **assignment problem**, which can be efficiently solved with an algorithm. But no closed formula exists for its solution.



Hence we know that the quantum problem can not, in general, be solved with a closed formula.

References:

- Ergotropy:

Allahverdjan et al. EPL 67 (2004)

- Passive states:

Lenard JSP 19 (1978), Alicki JPA (1979)

- Extraction of work from local systems:

Mukherjee et al. Pyhs. Rev. E 93 (2016)

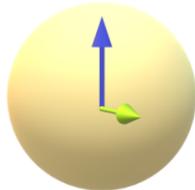
Alhambra et al. PRL 123 (2019)

...but it can be solved for qubits...

We derive an exact closed formula for the local ergotropy in the special case in which the local system S is a **two-level system**:

$$\mathcal{E}_S(\hat{\rho}_{SE}, \hat{H}_{SE}) = \begin{cases} \text{Tr}[|\mathcal{M}| - \mathcal{M}], & \text{for } \det[\mathcal{M}] \geq 0, \\ \text{Tr}[|\mathcal{M}| - \mathcal{M}] - \frac{2}{\|\mathcal{M}^{-1}\|}, & \text{for } \det[\mathcal{M}] < 0, \end{cases}$$

where M is a 3x3 matrix which depends in a complicated way on the state and on the Hamiltonian, and which is calculated using the formalism of **Generalized Pauli Operators**



...or relaxed to a SDP

The convexity of the local ergotropy functional allows us to relax the problem from a maximum over local unitaries to a maximum over **convex combination** of local unitaries. If we relax further to local **unital** transformations with the Choi isomorphism we can formulate a Semidefinite Programming bound:

$$\mathcal{E}_S(\hat{\rho}_{SE}, \hat{H}_{SE}) \leq \text{Tr}[\hat{H}_{SE} \hat{\rho}_{SE}] - \min_{E_{SS'}} \text{Tr}[C_{SS'} E_{SS'}],$$

with

$$\begin{aligned} E_{SS'} &\geq 0; \\ \text{Tr}_S E_{SS'} &= \mathbb{I}_{S'}; \\ \text{Tr}_{S'} E_{SS'} &= \mathbb{I}_S, \end{aligned}$$

Example: atom in a cavity

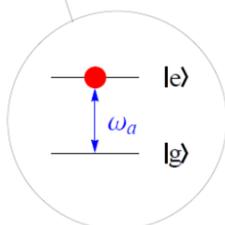
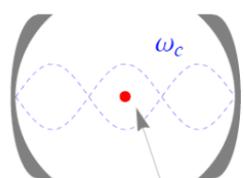
Jaynes-Cummings Hamiltonian:

$$\hat{H}_{SE}^{(JC)} := \hbar\omega_E \hat{a}^\dagger \hat{a} + \frac{\hbar\omega_S}{2} \hat{\sigma}_{S,3} + \frac{\hbar\Omega}{2} (\hat{a} \otimes \hat{\sigma}_S^+ + \hat{a}^\dagger \otimes \hat{\sigma}_S^-)$$

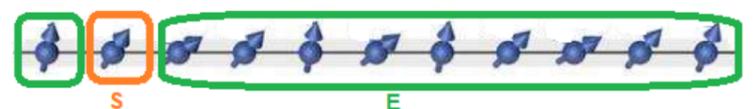
We find that the local ergotropy of the *dressed eigenstates* of an atom in a cavity is always bigger than the ergotropy of the isolated atom (taking into account the energetic cost of switching the interaction off).

Local ergotropy of the $|n, +\rangle$ eigenstate:

$$\mathcal{E}_S(|n, +\rangle, \hat{H}_{SE}^{(JC)}) = \hbar\omega_S \cos 2\theta_n + \frac{\sqrt{n+1}}{2} \hbar|\Omega| \sin 2\theta_n,$$



Example: XXZ spin chain



$$\mathcal{E}_S(\hat{\rho}_{SE}; \hat{H}_{SE}) = \begin{cases} 0 & (|k| \leq N/4) \\ -\frac{8J}{N} \cos\left(\frac{2\pi k}{N}\right) & (|k| > N/4) \end{cases}$$

The local ergotropy of a spin site in a XXZ chain can be nonzero even if the local state is passive (i.e. would have zero ergotropy)

