

ROLE OF COHERENCE AND DEGENERACIES IN QUANTUM SYNCHRONISATION

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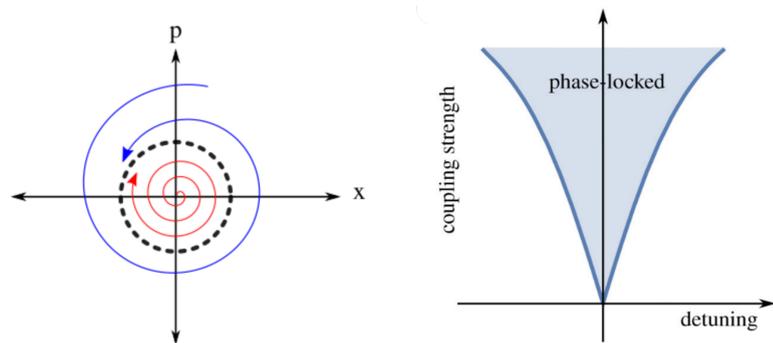
OBJECTIVES

- Coherent output of the heat engines/lasers have been related to the coherences present in the system.
- Coherences also implies the steady state synchronisation in quantum limit cycle oscillators.
- Recently, relationship between the power output and steady state synchronisation of Scovil and Schulz-Dubois heat engine was established[1] which we discuss here in detail.
- In our work, we investigate the change in relationship between the synchronisation and thermodynamic quantities for two mutually **coupled thermal machines**.

WHAT IS SYNCHRONISATION?

“Synchronisation is the adjustment of rhythms of weakly coupled self sustained oscillators”.

- Basic features of self sustained oscillators are:
 - 1 Dissipation, power source and non-linearity.
 - 2 Form of oscillation is independent of initial position.
 - 3 Oscillations are stable to (at least rather small) perturbations.



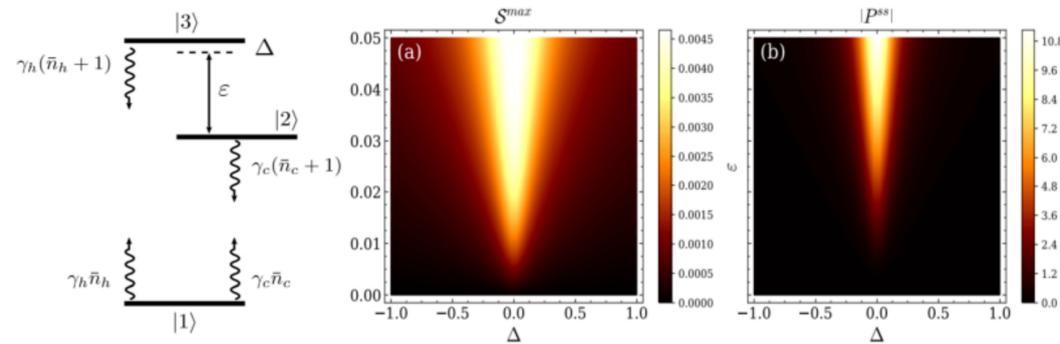
STEPS TO STUDY SYNC

- 1 Establish a valid **limit cycle**.
- 2 Apply a perturbation: **drive/coupling**.
- 3 A valid **synchronisation measure** to account for phase locking.

SYNC IN SCOVIL AND SCHULZ-DUBOIS HEAT ENGINE[1]

Establishing a valid **limit cycle**:

- Coherent state: $|n_3\rangle = |\theta, \zeta, \phi_1, \phi_2\rangle = |\cos(\theta), e^{i\phi_1}\cos\zeta\sin\theta, e^{i\phi_2}\sin\zeta\sin\theta\rangle$.
- $e^{iH_0t}|\theta, \zeta, \phi_1, \phi_2\rangle = |\theta, \zeta, \phi_1 - \omega_{21}t, \phi_2 - \omega_{31}t\rangle$; $H_0 = \sum_i \omega_i |i\rangle\langle i|$.
- ϕ_i 's are free phases which can be locked on by external drive/other oscillator.
- Master eqn.: $\dot{\rho}^R = -i[H_0^R + V, \rho^R] + \mathcal{D}_{13}^h[\rho] + \mathcal{D}_{12}^c[\rho^R]$; $H_0^R = \Delta\sigma_{33}$, $V = \epsilon S_x$.
- $Q(\theta, \zeta, \phi_1, \phi_2) = \langle n_3 | \rho | n_3 \rangle$ shows equiprobable distribution of phases ϕ_i 's for the steady state in absence of drive ($\epsilon = 0$) representing limit cycle behaviour.



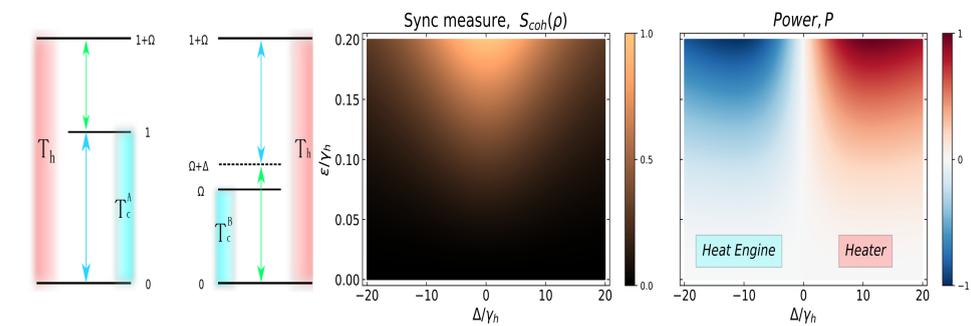
- **Perturbation**: External drive, $\epsilon \neq 0 \rightarrow$ phases ϕ_i 's can be localized.
- **Synchronisation measure**: $S(\phi_1, \phi_2) = \int d\Omega Q(\theta, \zeta, \phi_1, \phi_2) - \frac{1}{4\pi^2}$;
- $\dot{E} = \text{Tr}\{H_0^R \dot{\rho}^R\} = \underbrace{-i\text{Tr}\{[H_0^R + V, \rho^R]H_0^R\}}_{\text{Power}} + \underbrace{\text{Tr}\{\sum_i \mathcal{D}_i[\rho]H_0^R\}}_{\text{Heat}}$.
- Steady state power: $P_{ss} = -i\text{Tr}([H_0^R + V, \rho_{ss}]H_0^R) = 2\epsilon\omega_{32}\text{Im}[\rho_{23}^{ss}]$.

VALIDITY OF ‘LOCAL’ MASTER EQUATION

- ‘Global’ master equation produces thermodynamically consistent results but require a clean timescale separation which might break down in presence of degeneracies.
- ‘Local dissipators’ are found to coincides with the zeroth order of the Redfield dissipators if degeneracy of the bare Hamiltonian H_0 changes after turning on the interaction term V .

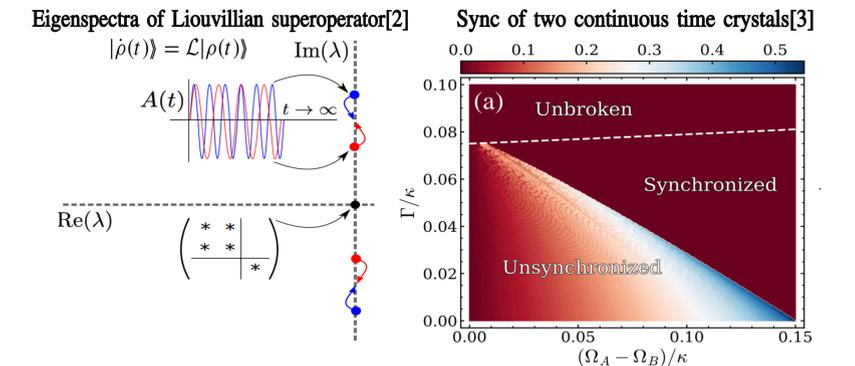
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TWO COUPLED THERMAL MACHINE!![2]



- $H_0 = \sigma_{22}^A + (1 + \Omega)\sigma_{33}^A + (\Omega + \Delta)\sigma_{22}^B + (1 + \Omega)\sigma_{33}^B$.
- $V = \epsilon(\sigma_{23}^A\sigma_{21}^B + \sigma_{12}^A\sigma_{32}^B + h.c.)$.
- For $\Delta = 0$ interaction is energy conserving, i.e. $[H_0, V = 0]$.
- Relative entropy of synchronisation: $S_{coh}(\rho) = S(\rho_{diag}) - S(\rho)$
- $P = -i\text{Tr}([H_0 + V, \rho]H_0) = 2g\Delta \text{Im}[\rho_{35}^{ss} + \rho_{75}^{ss}]$; $\rho^{ss} = \sum_{i,j=1}^9 \rho_{ij}^{ss} |i\rangle\langle j|$.

A DIFFERENT VARIETY OF SYNC[2, 3]



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- [1] Noufal Jaseem, Michal Hajdušek, Vlatko Vedral, Rosario Fazio, Leong-Chuan Kwek, and Sai Vinjanampathy. *Phys. Rev. E*, 101:020201, Feb 2020.
- [2] Parvinder Solanki, Noufal Jaseem, Michal Hajdušek, and Sai Vinjanampathy. *Physical Review A*, 105(2):L020401, 2022.
- [3] Michal Hajdušek, Parvinder Solanki, Rosario Fazio, and Sai Vinjanampathy. *Physical Review Letters*, 128(8):080603, 2022.